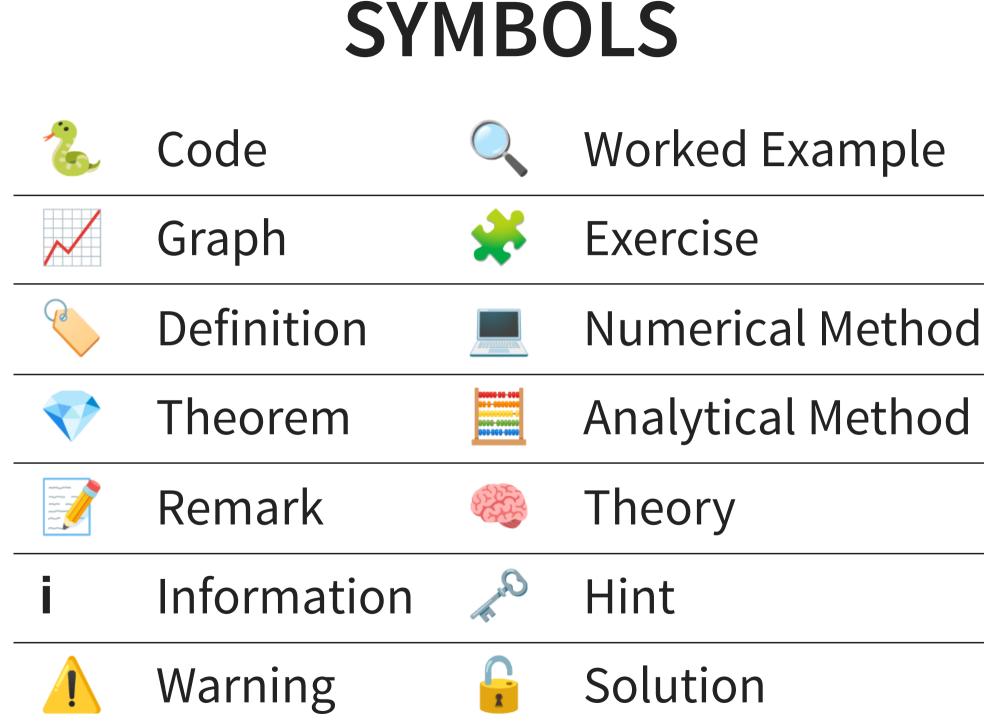
MODELS



CONTROL ENGINEERING WITH PYTHON

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from numpy import *
from numpy.linalg import *
from matplotlib.pyplot import *



ORDINARY DIFFERENTIAL EQUATION (ODE)

The "simple" version:

$$\dot{x} = f(x)$$

where:

- State: $x \in \mathbb{R}^n$
- State space: \mathbb{R}^n
- Vector field: $f : \mathbb{R}^n \to \mathbb{R}^n$.



More general versions:

• Time-dependent vector-field:

$$\dot{x}=f(t,x),\;t\in I\subset \mathbb{R},$$

- $x \in X$, open subset of \mathbb{R}^n ,
- $x \in X$, n-dimensional manifold.

VECTOR FIELD

- Visualize f(x) as an **arrow** with origin the **point** x.
- Visualize *f* as a field of such arrows.
- In the plane (n = 2), use quiver from Matplotlib.



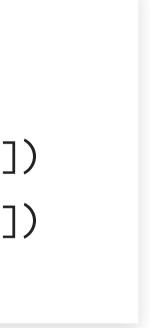
We define a Q function helper whose arguments are

- f: the vector field (a function)
- xs, ys: the coordinates (two 1d arrays)

and which returns:

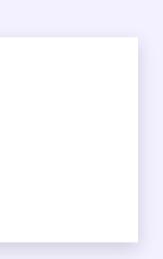
• the tuple of arguments expected by quiver.

def Q(f, xs, ys):
 X, Y = meshgrid(xs, ys)
 fx = vectorize(lambda x, y: f([x, y])[0])
 fy = vectorize(lambda x, y: f([x, y])[1])
 return X, Y, fx(X, Y), fy(X, Y)



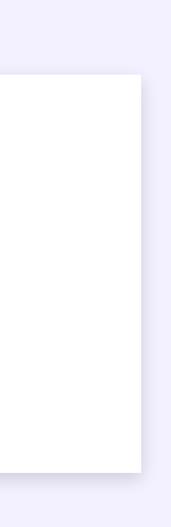
Consider f(x,y) = (-y,x).

def f(xy):
 x, y = xy
 return array([-y, x])





figure()
x = y = linspace(-1.0, 1.0, 20)
ticks = [-1.0, 0.0, 1.0]
xticks(ticks); yticks(ticks)
gca().set_aspect(1.0)
quiver(*Q(f, x, y))





A solution of $\dot{x} = f(x)$ is

- a (continuously) differentiable function $x: I
 ightarrow \mathbb{R}^n,$
- defined on a (possibly unbounded) interval I of \mathbb{R} ,
- such that for every $t\in I,$

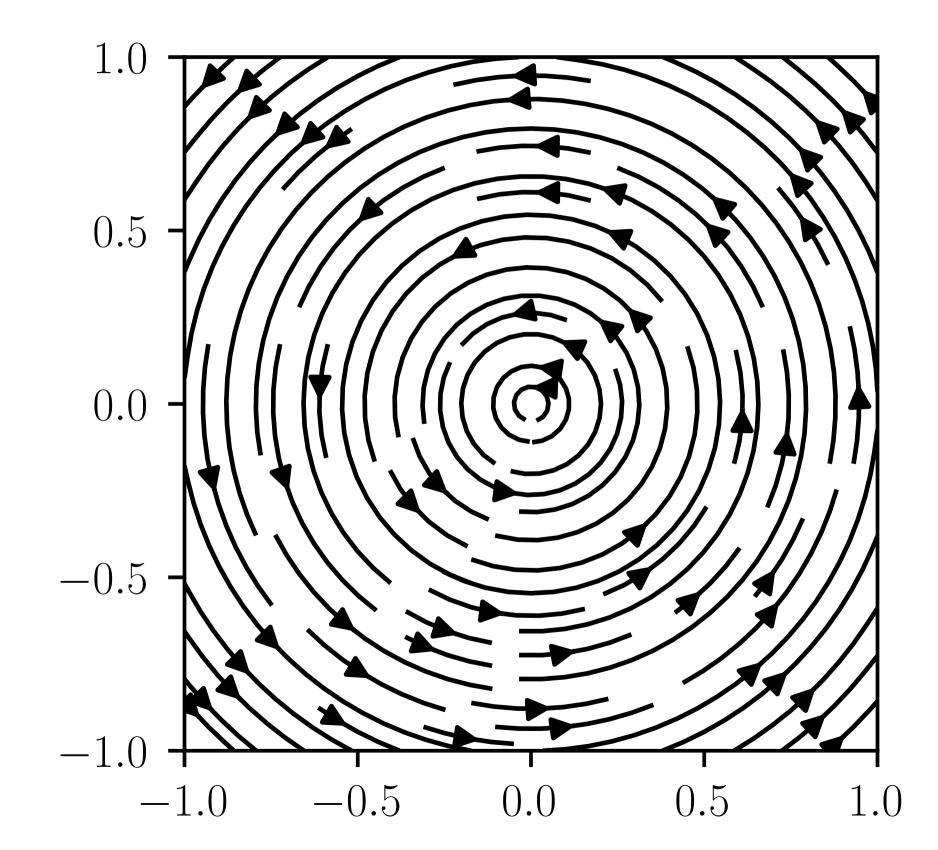
$$\dot{x}(t) = dx(t)/dt = f(x(t)).$$

$egin{array}{ll} I o {\mathbb R}^n, \ I ext{ of } {\mathbb R}, \end{array}$



When n = 2, represent a diverse set of solutions in the state space with streamplot

figure() x = y = linspace(-1.0, 1.0, 20)gca().set_aspect(1.0) streamplot(*Q(f, x, y), color="k")



INITIAL VALUE PROBLEM (IVP) Solutions x(t), for $t \geq t_0$, of

$$\dot{x} = f(x)$$

such that

$$x(t_0)=x_0\in \mathbb{R}^n.$$

The initial condition (t_0, x_0) is made of

• the initial time $t_0 \in \mathbb{R}$ and

• the initial value or initial state $x_0 \in \mathbb{R}^n$.

The point x(t) is the state at time t.



HIGHER-ORDER ODES

(Scalar) differential equations whose structure is

$$y^{(n)}(t)=g(y,\dot{y},\ddot{y},\ldots,y^{(n-1)})$$

where n > 1.



The previous *n*-th order ODE is equivalent to the firstorder ODE

$$\dot{x}=f(x),\,x\in\mathbb{R}^n$$

with

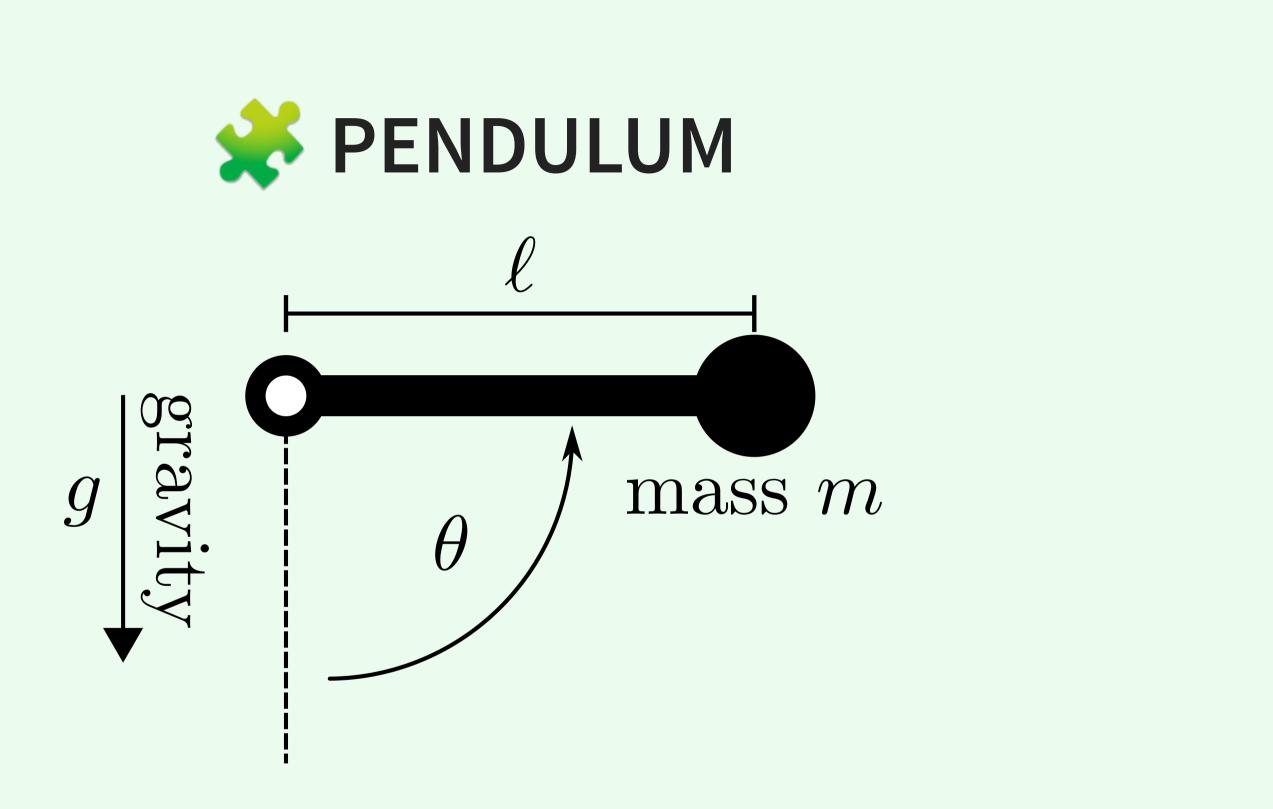
$$f(y_0,\ldots,y_{n-2},y_{n-1}):=(y_1,\ldots,y_{n-1},g(y_n))$$

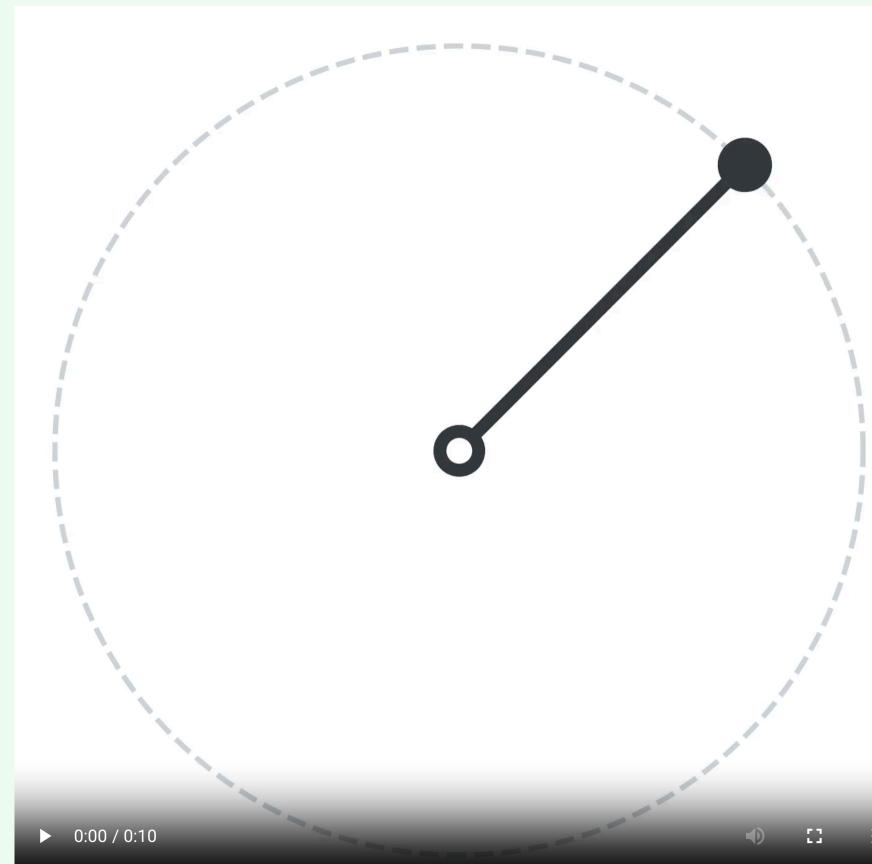
S he first-

$y_0,\ldots,y_{n-1})).$

The result is more obvious if we expand the first-order equation:

$$egin{array}{rll} \dot{y}_0 &=& y_1 \ \dot{y}_1 &=& y_2 \ dots &dots ˙$$







Establish the equations governing the pendulum dynamics.



Generalize the dynamics when there is a friction torque $c = -b\dot{\theta}$ for some $b \ge 0$.

We denote ω the pendulum **angular velocity**:

$$\omega := \dot{ heta}.$$



Transform the dynamics into a first-order ODE with state $x = (\theta, \omega)$.



Draw the system stream plot when $m=1, \ell=1,$ g = 9.81 and b = 0.



Determine least possible angular velocity $\omega_0 > 0$ such that when heta(0)=0 and $\dot{ heta}(0)=\omega_0$, the pendulum reaches (or overshoots) $\theta(t) = \pi$ for some t > 0.



1. 🔓

The pendulum total mechanical energy E is the sum of its kinetic energy K and its potential energy V:

$$E = K + V.$$

The kinetic energy depends on the mass velocity v:

$$K=rac{1}{2}mv^2=rac{1}{2}m\ell^2\dot{ heta}^2$$

The potential energy mass depends on the pendulum elevation y. If we set the reference y = 0 when the pendulum is horizontal, we have

$$V = mgy = -mg\ell\cos\theta$$

$$\Rightarrow ~E = K + V = rac{1}{2}m\ell^2\dot{ heta}^2 - mg\ell \cos^2 \phi^2$$

If the system evolves without any energy dissipation,

$$egin{aligned} \dot{E} &= rac{d}{dt}igg(rac{1}{2}m\ell^2\dot{ heta}^2 - mg\ell\cos hetaigg) \ &= m\ell^2\dot{ heta}\ddot{ heta} + mg\ell(\sin heta)\dot{ heta} \ &= 0 \end{aligned}$$

 $\Rightarrow m\ell^2\ddot{ heta} + mg\ell\sin heta = 0.$

$\cos \theta$.

2. 🔓

When there is an additional dissipative torque $c = -b\theta$, we have instead

$$\dot{E}=c\dot{ heta}=-b\dot{ heta}^2$$

and thus

 $m\ell^2\ddot{ heta} + b\dot{ heta} + mg\ell\sin heta = 0.$

3. $\widehat{\theta}$. With $\omega := \dot{\theta}$, the dynamics becomes

$$egin{array}{rcl} \dot{ heta} &=& \omega \ \dot{\omega} &=& -(b/m\ell^2)\omega - (g/\ell)\sin heta \end{array}$$

4. 🔓

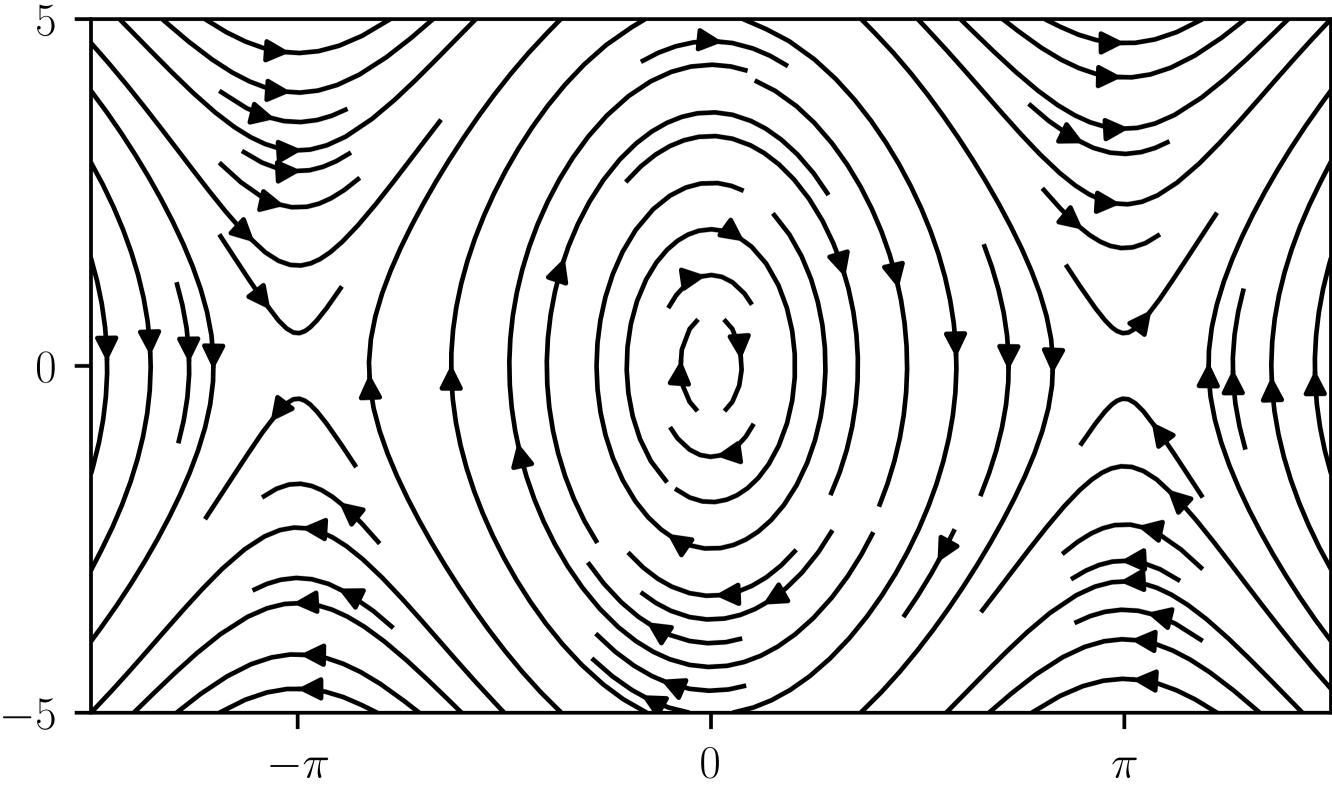
m=1.0; b=0.0; l=1.0; g=9.81
def f(theta_d_theta):
 theta, d_theta = theta_d_theta
 J = m * l * l
 d2_theta = - b / J * d_theta
 d2_theta += - g / l * sin(theta)
 return array([d_theta, d2_theta])





figure()
theta = linspace(-1.5 * pi, 1.5 * pi, 100)
d_theta = linspace(-5.0, 5.0, 100)
labels = [r"\$-\pi\$", "\$0\$", r"\$\pi\$"]
xticks([-pi, 0, pi], labels)
yticks([-5, 0, 5])
streamplot(*Q(f, theta, d_theta), color="k")





5. 🔓

In the top vertical configuration, the total mechanical energy of the pendulum is

$$E_{ op} = rac{1}{2} m \ell^2 \dot{ heta}^2 - mg\ell \cos \pi = rac{1}{2} m \ell^2 \dot{ heta}^2 + rac{1}{2} m \ell^2 \dot{ heta}^2$$
 .

Hence we have at least $E_{\perp} \geq mg\ell$.

 $+ mg\ell.$

On the other hand, in the bottom configuration,

$$E_{\perp} = rac{1}{2} m \ell^2 \dot{ heta}^2 - mg\ell \cos 0 = rac{1}{2} m \ell^2 \dot{ heta}^2 + rac{1}{2} m \ell^2 \dot{ heta}^2$$
 -

Hence, without any loss of energy, the initial velocity must satisfy $E_{\perp} \geq E_{\perp}$ for the mass to reach the top position.



 $-mg\ell.$

That is

$$E_{\perp} = rac{1}{2} m \ell^2 \dot{ heta}^2 - mg \ell \geq mg \ell = E$$

which leads to:

$$|\dot{ heta}| \geq 2\sqrt{rac{g}{\ell}}.$$

 $\mathcal{I}_{ op}$